# Physics 122: Electricity \& Magnetism - Lecture 8 Electric Potential 

Baris EMRE

## Potential due to a group of point charges

- Use superposition

$$
V=-\int_{\infty}^{r} \vec{E} \cdot d \vec{s}=-\sum_{i=1}^{n} \int_{\infty}^{r} \vec{E}_{i} \cdot d \vec{s}=\sum_{i=1}^{n} V_{i}
$$

- For point charges

$$
V=\sum_{i=1}^{n} V_{i}=\frac{1}{4 \pi \varepsilon_{0}} \sum_{i=1}^{n} \frac{q_{i}}{r_{i}}
$$

$\square$ The sum is an algebraic sum, not a vector sum.
$\square$ E may be zero where V does not equal to zero.
$\square \quad \mathrm{V}$ may be zero where E does not equal to zero.


## Electric Field and Electric Potential

4. Which of the following figures have $\mathrm{V}=0$ and

## $\mathrm{E}=0$ at red point?



A


C


D


B


## Potential due to a Continuous Charge Distribution

$\square$ Find an expression for dq:

- $d q=\lambda d /$ for a line distribution
- $d q=\sigma d A$ for a surface distribution
- $d q=\rho d V$ for a volume distribution
$\square$ Represent field contributions at P due to point charges dq located in the distribution.

$$
d V=\frac{1}{4 \pi \varepsilon_{0}} \frac{d q}{r}
$$

- Integrate the contributions over the whole distribution, varying the displacement as needed,

$$
V=\int d V=\frac{1}{4 \pi \varepsilon_{0}} \int \frac{d q}{r}
$$

## Example: Potential Due to a Charged Rod

- A rod of length $L$ located along the $x$ axis has a uniform linear charge density $\lambda$. Find the electric potential at a point $P$ located on the $y$ axis a distance $d$ from the origin.
- Start with $d q=\lambda d x$

$$
d V=\frac{1}{4 \pi \varepsilon_{0}} \frac{d q}{r}=\frac{1}{4 \pi \varepsilon_{0}} \frac{\lambda d x}{\left(x^{2}+d^{2}\right)^{1 / 2}}
$$

- then,

$$
\begin{aligned}
& V=\int d V=\int_{0}^{L} \frac{\lambda}{4 \pi \varepsilon_{0}} \frac{d x}{\left(x^{2}+d^{2}\right)^{1 / 2}}=\frac{\lambda}{4 \pi \varepsilon_{0}}\left[\ln \left(x+\left(x^{2}+d^{2}\right)^{1 / 2}\right)\right]_{0}^{L} \\
& =\frac{\lambda}{4 \pi \varepsilon_{0}}\left[\ln \left(L+\left(L^{2}+d^{2}\right)^{1 / 2}\right)-\ln d\right]
\end{aligned}
$$

$$
V=\frac{\lambda}{4 \pi \varepsilon_{0}} \ln \left[\frac{L+\left(L^{2}+d^{2}\right)^{1 / 2}}{d}\right]
$$


(a)

(b)

## Potential Due to a Charged Isolated Conductor

$\square$ According to Gauss' law, the charge resides on the conductor's outer surface.

- Furthermore, the electric field just outside the conductor is perpendicular to the surface and field inside is zero.
- Since

$$
V_{B}-V_{A}=-\int_{A}^{B} \vec{E} \cdot d \vec{s}=0
$$

- Every point on the surface of a charged conductor in equilibrium is at the same electric potential.
- Furthermore, the electric potential is constant everywhere inside the conductor and equal to its value to its value at the surface.
$V_{B}-V_{A}=-\int_{A}^{B} \vec{E} \cdot d \vec{s}=0$
(1)



## Calculating the Field from the Potential

- Suppose that a positive test charge $\mathrm{q}_{0}$ moves through a displacement ds from on equipotential surface to the adjacent surface.
- The work done by the electric field on the test charge is $W=-d U=-q_{0} d V$.
- The work done by the electric field may also be written as $W=q_{0} \vec{E} \cdot d \vec{s}$
- Then, we have

$$
-q_{0} d V=q_{0} E(\cos \theta) d s \quad E \cos \theta=-\frac{d V}{d s}
$$

$\square$ So, the component of E in any direction is the negative of the rate at which the electric potential changes with distance in that direction.

$$
E_{s}=-\frac{\partial V}{\partial s}
$$

- If we know $\mathrm{V}(\mathrm{x}, \mathrm{y}, \mathrm{z})$,

$$
E_{x}=-\frac{\partial V}{\partial x} \quad E_{y}=-\frac{\partial V}{\partial y} \quad E_{z}=-\frac{\partial V}{\partial z}
$$

## Electric Potential Energy of a System of Point Charges

$$
\begin{array}{cc}
\Delta U=U_{f}-U_{i}=-W & W=\vec{F} \cdot \Delta \vec{r}=q \vec{E} \cdot \Delta \vec{r} \\
W_{a p p}=-W & \Delta U=U_{f}-U_{i}=W_{a p p}
\end{array}
$$

- Start with (set $\mathrm{U}_{\mathrm{i}}=0$ at $\infty$ and $\mathrm{U}_{\mathrm{f}}=\mathrm{U}$ at r )

$$
V=\frac{1}{4 \pi \varepsilon_{0}} \frac{q_{1}}{r}
$$

- We have

$$
U=q_{2} V=\frac{1}{4 \pi \varepsilon_{0}} \frac{q_{1} q_{2}}{r}
$$

- If the system consists of more than two charged particles, calculate $U$ for each pair of charges and sum the terms algebraically.
$U=U_{12}+U_{13}+U_{23}=\frac{1}{4 \pi \varepsilon_{0}}\left(\frac{q_{1} q_{2}}{r_{12}}+\frac{q_{1} q_{3}}{r_{13}}+\frac{q_{2} q_{3}}{r_{23}}\right)$


## Summary

$\square$ Electric Potential Energy: a point charge moves from $i$ to $f$ in an electric field, the change in electric potential energy is
$\square$ Electric Potential Difference between two points $i$ and $f$ in an electric field:

- Equipotential surface: the points on it all have the same electric potential. No work is done while moving charge on it. The electric field is always directed perpendicularly to corresponding equipotential surfaces.
$\square$ Finding V from E:
$\square$ Potential due to point charges: $\quad \Delta V \equiv \frac{\Delta U}{q_{0}}=-\int_{i}^{\mu} \vec{E} \cdot d \vec{s}$
$\square$ Potential due to a collection of point charges:

$$
V=\sum_{i=1}^{n} V_{i}=\frac{1}{4 \pi \varepsilon_{0}} \sum_{i=1}^{n} \frac{q_{i}}{r_{i}}
$$

$\square$ Potential due to a continuous charge distribution:
$\square$ Potential of a charged conductor is constant everywhere inside the conductor and equal to its value to its value at the surface.

- Calculatiing E from V: $\quad E_{s}=-\frac{\partial V}{\partial s} \quad E_{x}=-\frac{\partial V}{\partial x} \quad E_{y}=-\frac{\partial V}{\partial y} \quad E_{z}=-\frac{\partial V}{\partial z}$
$\square$ Electric potential energy of system of point charges:

$$
\Delta U=U_{f}-U_{i}=-W
$$

$$
\Delta V=V_{f}-V_{i}=\frac{U_{f}}{q}-\frac{U_{i}}{q}=\frac{\Delta U}{q}
$$

$$
V(r)=\frac{1}{4 \pi \varepsilon_{0}} \frac{q}{r}
$$

$$
V=\int d V=\frac{1}{4 \pi \varepsilon_{0}} \int \frac{d q}{r}
$$

$$
U=q_{2} V=\frac{1}{4 \pi \varepsilon_{0}} \frac{q_{1} q_{2}}{r}
$$

