

Joule's law

In the presence of an electric field \vec{E} , free electrons in a conductor have a drift (average) velocity, \vec{v} . Collisions among free electrons and immobile atoms transfer energy from the electric field to thermal vibration. Quantitatively, the work done by \vec{E} in moving an amount of charge q in unit volume for a differential “drift” displacement $\vec{dl} = \vec{v}dt$ is

	$\Delta W = \vec{F} \cdot \vec{dl} = q\vec{E} \cdot \vec{v}\Delta t$	
	$\Delta W = q\vec{E} \cdot \vec{dl} = q\vec{E} \cdot \vec{v}\Delta t$	
Power dissipated is,		
	$p = \frac{\Delta W}{\Delta t} = \frac{q\vec{E} \cdot \vec{v}\Delta t}{\Delta t}$	
	$p = \frac{q\vec{E} \cdot \vec{v}\Delta t}{\Delta t}$	
	$p = q\vec{E} \cdot \vec{v}$	
	$p = q\vec{E} \cdot \frac{\vec{J}}{q}$	
	$p = \vec{E} \cdot \vec{J}$	Per unit volume
(Ohmic) Power Density		
	$p = \vec{E} \cdot \vec{J}$	(Watts/m ³)
Total Power Dissipation with inhomogeneous $\vec{E}(R)$ & $\sigma(R), \vec{J}(R)$		

$$P = \int p dV$$

$$P = \int_V \vec{E}(R) \cdot \vec{J}(R) dV \quad (\text{Watts})$$

If we apply a voltage difference V_{12} across a homogeneous conductor of uniform cross-sectional area S , conductivity σ and length L , \Rightarrow

$$P = \int_V \vec{E}(R) \cdot \vec{J}(R) dV = \int_C \vec{E}(R) \cdot \vec{J}(R) S dl =$$

$$P = \int_C \vec{E}(R) \cdot J(R) \hat{l} \cdot S \cdot dl$$

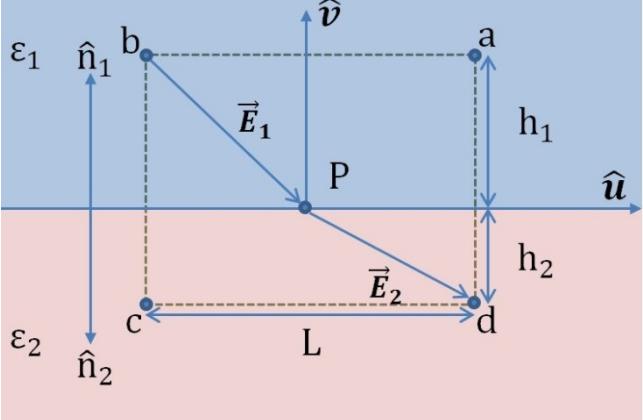
$$P = \int_C \vec{E}(R) \cdot I \hat{l} dl = \int_C E(R) \cdot I \cdot dl =$$

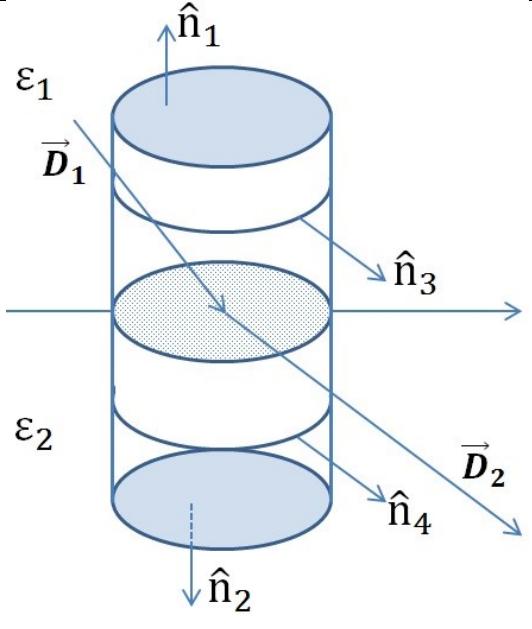
$$= \int_C \vec{E}(R) \cdot I \hat{l} dl = I \int_C E(R) dl = IV_{12}$$

Boundary Conditions

$$\nabla \times \vec{E} = \vec{0}$$

$$\int_S \nabla \times \vec{E} \cdot \vec{dS} = \oint_C \vec{E} \cdot \vec{dl} = 0$$

	
	$E_{1//} = E_{2//}$
	Tangential component of \vec{E}_1 is equal to tangential component of \vec{E}_2
=>	$\frac{J_{1//}}{\sigma_1} = \frac{J_{2//}}{\sigma_2}$

Boundary Condition 2	
	$\nabla \cdot \vec{J} = 0$
	$\int_V \nabla \cdot \vec{J} dV = \int_V \mathbf{0} dV$
	$\oint_S \vec{J} \cdot \vec{dS} = \int_V \mathbf{0} dV$ $\oint_S \vec{J} \cdot \vec{dS} = 0$

$$\oint_S \vec{J} \cdot \overrightarrow{dS} = \int_{S_1} \vec{J}_1 \cdot \overrightarrow{dS} + \int_{S_2} \vec{J}_2 \cdot \overrightarrow{dS} + \int_{S_3} \vec{J}_1 \cdot \overrightarrow{dS} + \int_{S_4} \vec{J}_2 \cdot \overrightarrow{dS}$$

$$\int_{S_3} \vec{J}_1 \cdot \overrightarrow{dS} \rightarrow \mathbf{0}$$

$$\int_{S_4} \vec{J}_2 \cdot \overrightarrow{dS} \rightarrow \mathbf{0}$$

$$\oint_S \vec{J} \cdot \overrightarrow{dS} = \mathbf{0} = \int_{S_1} \vec{J}_1 \cdot \overrightarrow{dS} + \int_{S_3} \vec{J}_3 \cdot \overrightarrow{dS}$$

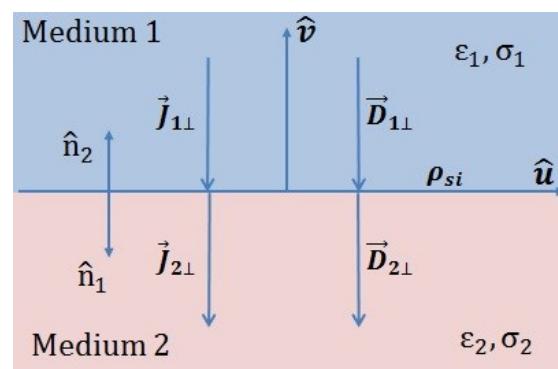
$$J_{1\perp}(-S) + J_{2\perp}(+S) = \mathbf{0}$$

$$\Rightarrow J_{2\perp} - J_{1\perp} = \mathbf{0}$$

$$\Rightarrow \hat{n}_1 \cdot (\vec{J}_2 - \vec{J}_1) = \mathbf{0}$$

$$\Rightarrow \hat{n}_2 \cdot (\vec{J}_1 - \vec{J}_2) = \mathbf{0}$$

Ex: Find ρ_{si} in terms of $D_{1\perp}, D_{2\perp}, \varepsilon_1, \varepsilon_2, \sigma_1, \sigma_2$



$$\vec{D}_{1\perp} = D_{1\perp}(-\hat{v})$$

$$\vec{D}_{2\perp} = D_{2\perp}(-\hat{v})$$

$$D_{2\perp} - D_{1\perp} = \rho_{Si}$$

$$\epsilon_2 E_{2\perp} - \epsilon_1 E_{1\perp} = \rho_{Si}$$

$$J_{2\perp} - J_{1\perp} = \mathbf{0}$$

$$\sigma_2 E_{2\perp} - \sigma_1 E_{1\perp} = \mathbf{0}$$

$$E_{1\perp} = \frac{\sigma_2}{\sigma_1} E_{2\perp}$$

$$D_{1\perp} = \epsilon_1 E_{1\perp} = \epsilon_1 \frac{\sigma_2}{\sigma_1} E_{2\perp}$$

$$D_{1\perp} = \frac{\epsilon_1 \sigma_2}{\epsilon_2 \sigma_1} D_{2\perp}$$

$$E_{2\perp} = \frac{\sigma_1}{\sigma_2} E_{1\perp}$$

$$D_{2\perp} = \epsilon_2 E_{2\perp} = \epsilon_2 \frac{\sigma_1}{\sigma_2} E_{1\perp}$$

$$D_{2\perp} = \frac{\epsilon_2 \sigma_1}{\epsilon_1 \sigma_2} D_{1\perp}$$

$$\rho_{Si} = (1 - \frac{\epsilon_1 \sigma_2}{\epsilon_2 \sigma_1}) D_{2\perp}$$

$$\rho_{Si} = (\frac{\epsilon_2 \sigma_1}{\epsilon_1 \sigma_2} D_{1\perp} - 1) D_{1\perp}$$

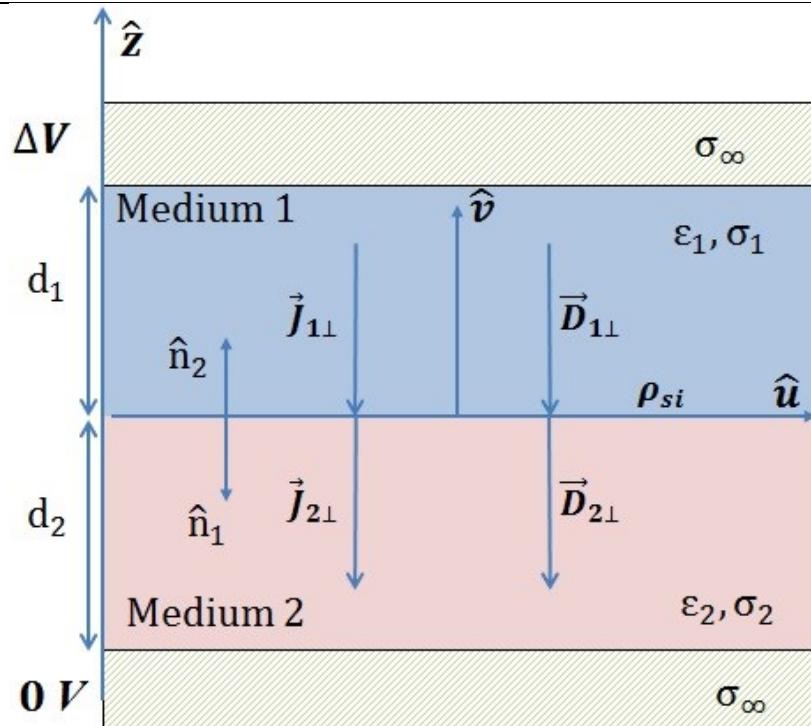
$$\rho_{Si} = \mathbf{0} \text{ if } \frac{\epsilon_2 \sigma_1}{\epsilon_1 \sigma_2} = 1, (\frac{\sigma_1}{\epsilon_1} = \frac{\sigma_2}{\epsilon_2})$$

Or,

$$\rho_{Si} = \mathbf{0} \text{ if } \sigma_1 = \sigma_2 = \mathbf{0}$$

Both media is lossless, thus no free charge can exist

Ex: Find J , E in the two lossy media between two parallel conducting plates biased by a dc voltage V_0 . Also find the surface charge densities on the two conducting plates and on the interface between the two lossy media.



$$\vec{J}_1 = J_{1\perp}(-\hat{\mathbf{z}})$$

$$\vec{J}_1 = J_{2\perp}(-\hat{\mathbf{z}})$$

$$J_{1\perp} = J_{2\perp} = J$$

$$\vec{E}_1 = E_1(-\hat{\mathbf{z}}) = E_{1\perp}(-\hat{\mathbf{z}})$$

$$\vec{E}_1 = E_1(-\hat{\mathbf{z}}) = \frac{J_{1\perp}}{\sigma_1}(-\hat{\mathbf{z}}) = \frac{J}{\sigma_1}(-\hat{\mathbf{z}})$$

$$\vec{E}_2 = E_2(-\hat{\mathbf{z}}) = E_{2\perp}(-\hat{\mathbf{z}})$$

$$\vec{E}_2 = \frac{J_{2\perp}}{\sigma_2}(-\hat{\mathbf{z}}) = \frac{J}{\sigma_2}(-\hat{\mathbf{z}})$$

$$\Delta V - 0 = - \int_{-}^{+} \vec{E} \cdot d\vec{l}$$

$$\Delta V = \left(- \int_0^{d_2} \vec{E}_2 \cdot d\vec{l}_2 \right) + \left(- \int_{d_2}^{d_2+d_1} \vec{E}_1 \cdot d\vec{l}_1 \right)$$

$$\Delta V = \left(- \int_0^{d_2} \vec{E}_2 \cdot dz\hat{\mathbf{z}} \right) + \left(- \int_{d_2}^{d_2+d_1} \vec{E}_1 \cdot dz\hat{\mathbf{z}} \right)$$

$$\Delta V = \left(- \int_0^{d_2} \frac{J}{\sigma_2}(-\hat{\mathbf{z}}) \cdot dz\hat{\mathbf{z}} \right) + \left(- \int_{d_2}^{d_2+d_1} \frac{J}{\sigma_1}(-\hat{\mathbf{z}}) \cdot dz\hat{\mathbf{z}} \right)$$

$$\Delta V = \left(\frac{J}{\sigma_2} \int_0^{d_2} (+\hat{\mathbf{z}}) \cdot dz\hat{\mathbf{z}} \right) + \left(\frac{J}{\sigma_1} \int_{d_2}^{d_2+d_1} (+\hat{\mathbf{z}}) \cdot dz\hat{\mathbf{z}} \right)$$

$$\Delta V = \left(\frac{J}{\sigma_2} \int_0^{d_2} dz \right) + \left(\frac{J}{\sigma_1} \int_{d_2}^{d_2+d_1} dz \right)$$

$$\Delta V = \left(\frac{J}{\sigma_2} d_2 \right) + \left(\frac{J}{\sigma_1} d_1 \right)$$

$$J = \frac{\Delta V}{\left(\frac{d_1}{\sigma_1} + \frac{d_2}{\sigma_2} \right)} = \frac{\sigma_1 \sigma_2 \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$$E_1 = \frac{J}{\sigma_1} = \frac{\sigma_2 \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$$E_2 = \frac{J}{\sigma_2} = \frac{\sigma_1 \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$$D_1 = \varepsilon_1 E_1 = \frac{\varepsilon_1 \sigma_2 \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$$D_2 = \varepsilon_2 E_2 = \frac{\varepsilon_2 \sigma_1 \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$$\hat{n}_{\sigma 1} \cdot (\vec{D}_{\varepsilon_1} - \vec{D}_{\sigma 2}) = \rho_{s1}$$

$$-\hat{\mathbf{z}} \cdot (\vec{D}_1 - \mathbf{0}) = -\hat{\mathbf{z}} \cdot (D_1(-\hat{\mathbf{z}}) - \mathbf{0}) = \rho_{s1}$$

$$D_1 = \rho_{s1}$$

$$\rho_{s1} = D_1 = \frac{\varepsilon_1 \sigma_2 \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$$\hat{n}_{\sigma 2} \cdot (\vec{D}_{\varepsilon_2} - \vec{D}_{\sigma 2}) = \rho_{s2}$$

$$+\hat{\mathbf{z}} \cdot (\vec{D}_2 - \mathbf{0}) = +\hat{\mathbf{z}} \cdot (D_2(-\hat{\mathbf{z}}) - \mathbf{0}) = \rho_{s2}$$

$$D_2 = -\rho_{s2}$$

$$\rho_{s2} = -D_2 = -\frac{\varepsilon_2 \sigma_1 \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$$\rho_{si} = (1 - \frac{\varepsilon_1 \sigma_2}{\varepsilon_2 \sigma_1}) D_{2\perp}$$

$$\rho_{si} = (1 - \frac{\varepsilon_1 \sigma_2}{\varepsilon_2 \sigma_1}) \frac{\varepsilon_2 \sigma_1 \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$$\rho_{si} = \frac{(\varepsilon_2 \sigma_1 - \varepsilon_1 \sigma_2) \Delta V}{(\sigma_2 d_1 + \sigma_1 d_2)}$$

$ \rho_{s1} \neq \rho_{s2} $ but $\rho_{s1} + \rho_{s2} + \rho_{si} = 0$	
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Evaluation of Resistance

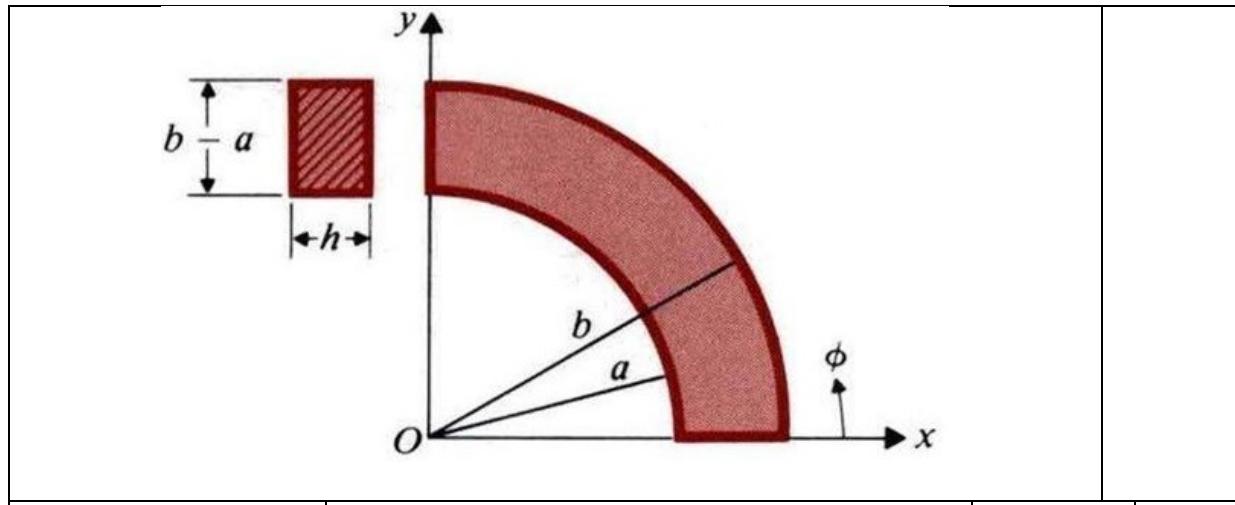
Resistance of single imperfect conductor

The resistance R of a piece of homogeneous lossy medium of finite conductivity σ can be evaluated by the following steps:

- (1) Assume a potential difference ΔV between the two conductors.
- (2). Find the potential distribution Φ by solving boundary-value problem.
- (3) Find \vec{E} by using $\vec{E} = \nabla\Phi$
- (4) Find the total current by $I = \int \vec{J} \cdot d\vec{S} = \int \sigma \vec{E} \cdot d\vec{S}$
- (5) $R = \frac{\Delta V}{I}$

Ex: Consider a quarter-circular washer of rectangular cross section and finite conductivity σ .

Find the resistance if the two electrodes are located at $\phi = 0$ and $\phi = \pi/2$.



BC_1	$\Phi(0) = \Delta V$		
BC_2	$\Phi\left(\frac{\pi}{2}\right) = 0$		

	$\Phi = \Phi(\emptyset)$		
No free charge	$\nabla \cdot \vec{D} = \rho_v = 0$		
	$\nabla \cdot \epsilon \vec{E} = 0$		
Simple medium	$\epsilon \nabla \cdot \vec{E} = 0$		
$-\nabla \Phi = \vec{E}$	$\epsilon \nabla \cdot (-\nabla \Phi) = 0$		
$\nabla \cdot \nabla = \nabla^2$	$-\epsilon \nabla \cdot \nabla \Phi = -\epsilon \nabla^2 \Phi = 0$		
	$\nabla^2 \Phi = 0$		
	$\frac{\partial^2 \Phi}{\partial \emptyset^2} = 0$		
	$\frac{\partial \Phi}{\partial \emptyset} = C_1$		
	$\Phi = \emptyset C_1 + C_2$		
Using BCs.,	$\Phi(0) = \Delta V = 0 C_1 + C_2$		
	$\Phi\left(\frac{\pi}{2}\right) = 0 = \frac{\pi}{2} C_1 + C_2$		
	$C_1 = \frac{2\Delta V}{\pi}$		
	$C_2 = 0$		
	$\Phi = \frac{2\Delta V}{\pi} \emptyset$		
	$\vec{E} = -\nabla \Phi = -\hat{\emptyset} \frac{2\Delta V \mathbf{1}}{\pi r} = -\frac{2V_0 \mathbf{1}}{\pi r} \hat{\emptyset}$		
	$\vec{E} = \vec{E}(r) \hat{\emptyset}$		

$$\vec{J} = \sigma \vec{E} = -\sigma \frac{2\Delta V}{\pi} \frac{\mathbf{1}}{r} \hat{\phi}$$

$$I = \int \vec{J} \cdot \overrightarrow{dS}$$

$$I = \iint_{0 \ a}^{h \ b} -\sigma \frac{2\Delta V}{\pi} \frac{x \mathbf{d} \mathbf{1}}{r} \hat{\phi} \cdot (-\hat{\phi}) dr dz$$

$$= \left(\int_a^b +\sigma \frac{2\Delta V}{\pi} \frac{\mathbf{1}}{r} r dr \right) \left(\int_0^h dz \right)$$

$$= \left(\sigma \frac{2\Delta V}{\pi} \right) \left(\int_a^b \frac{\mathbf{1}}{r} r dr \right) \left(\int_0^h dz \right)$$

$$I = \left(\sigma \frac{2\Delta V}{\pi} \right) \cdot \ln \left(\frac{b}{a} \right) \cdot h$$

$$R = \frac{\Delta V}{I}$$

$$R = \frac{\Delta V}{\left(\sigma \frac{2\Delta V}{\pi} \right) \cdot \ln \left(\frac{b}{a} \right) \cdot h}$$

$$R = \left(\left(\sigma \frac{2\Delta V}{\pi} \right) \cdot \ln \left(\frac{b}{a} \right) \cdot h \right)^{-1} \quad (\text{Ohm})$$