Again from slide 13,

$$\vec{H}_{\perp i} = \left(\frac{E_{\perp i}}{\eta_1}\right) \left(e^{-j\beta_1 (x\sin\theta_i + z\cos\theta_i)}\right) \left(-\cos\theta_i \,\hat{x} + \,\sin\theta_i \,\hat{z}\right)$$

$$\vec{H}_{\perp r} = \left(\frac{E_{\perp r}}{\eta_1}\right) \left(e^{-j\beta_1 (x\sin\theta_r - z\cos\theta_r)}\right) \left(\cos\theta_r \,\hat{x} + \,\sin\theta_r \,\hat{z}\right)$$

$$\vec{H}_{\perp t} = (\frac{E_{\perp t}}{\eta_2}) (e^{-j\beta_2(x\sin\theta_t + z\cos\theta_t)}) (-\cos\theta_t \,\hat{x} + \frac{\sin\theta_t \,\hat{z}}{\cos\theta_t \,\hat{x}})$$
Lossless H_{\perp}

New At the boundary (z=0 / on xy-plane), the tangential component (x- and y-component) is

$$\vec{H}_{\perp r,tangential} = \left(\frac{E_{\perp r}}{\eta_1}\right) \left(e^{-j\beta_1 (x\sin\theta_r - 0\cos\theta_r)}\right) \left(\cos\theta_r \,\hat{x}\right)$$

$$\vec{H}_{\perp r,tangential} = \left(\frac{E_{\perp r}}{\eta_1}\right) \left(e^{-j\beta_1(x\sin\theta_r)}\right) \left(\cos\theta_r \,\hat{x}\right)$$

$$z=0 \text{ (on the boundary surface, xy-plane), Lossless, } H_{\perp} |_{21}$$

Again from slide 13,

$$\vec{H}_{\perp i} = \left(\frac{E_{\perp i}}{\eta_1}\right) \left(e^{-j\beta_1 (x\sin\theta_i + z\cos\theta_i)}\right) \left(-\cos\theta_i \,\hat{x} + \,\sin\theta_i \,\hat{z}\right)$$

$$\vec{H}_{\perp r} = \left(\frac{E_{\perp r}}{\eta_1}\right) \left(e^{-j\beta_1 (x\sin\theta_r - z\cos\theta_r)}\right) \left(\cos\theta_r \,\hat{x} + \,\sin\theta_r \,\hat{z}\right)$$

$$\vec{H}_{\perp t} = (\frac{E_{\perp t}}{\eta_2}) (e^{-j\beta_2(x\sin\theta_t + z\cos\theta_t)}) (-\cos\theta_t \,\hat{x} + \frac{\sin\theta_t \,\hat{z}}{\cos\theta_t \,\hat{x}})$$
Lossless H_{\perp}

New At the boundary (*z*=0 / on *xy*-plane), the tangential component (x- and y-component) is

$$\vec{H}_{\perp t, tangential} = \left(\frac{E_{\perp t}}{\eta_2}\right) \left(e^{-j\beta_2 (x\sin\theta_t + 0\cos\theta_r)}\right) \left(-\cos\theta_t \,\hat{x}\right)$$

$$\vec{H}_{\perp t,tangential} = \left(\frac{E_{\perp t}}{\eta_2}\right) \left(e^{-j\beta_2(x\sin\theta_t)}\right) \left(-\cos\theta_t \hat{x}\right)$$

$$z=0 \text{ (on the boundary stirface, xy-plane), Lossless, } H_{\perp}$$

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Summary of the tangential components of the magnetic field vectors at the boundary which is the *xy*-plane or equivalently *z*=0

$$\vec{H}_{\perp i,tangential} = \left(\frac{E_{\perp i}}{\eta_1}\right) \left(e^{-j\beta_1(x\sin\theta_i)}\right) \left(-\cos\theta_i \,\hat{x}\right)$$
$$\vec{H}_{\perp r,tangential} = \left(\frac{E_{\perp r}}{\eta_1}\right) \left(e^{-j\beta_1(x\sin\theta_r)}\right) \left(\cos\theta_r \,\hat{x}\right)$$
$$\vec{H}_{\perp t,tangential} = \left(\frac{E_{\perp t}}{\eta_2}\right) \left(e^{-j\beta_2(x\sin\theta_t)}\right) \left(-\cos\theta_t \,\hat{x}\right)$$
$$z=0 \text{ (on the boundary surface, xy-plane), Lossless, } H_{\perp}$$

In general,

$$\vec{H}_{\perp 1} = \vec{H}_{\perp i} + \vec{H}_{\perp r}$$
$$\vec{H}_{\perp 2} = \vec{H}_{\perp t}$$

Boundary Condition

$$\hat{n}_{2} \times \left(\hat{H}_{1} - \hat{H}_{2} \right) = \vec{J}_{S}$$

Boundary Condition for lossless case:

 $\vec{J}_{S} = 0$ (No conduction current at the boundary)

$$\vec{H}_{\perp 1,tangential} = \vec{H}_{\perp 2,tangential}$$

Using BC at the boundary (z=0)

$$\vec{H}_{\perp i,tangential} + \vec{H}_{\perp r,tangential} = \vec{H}_{\perp t,tangential}$$

$$\left(\frac{E_{\perp i}}{\eta_1}\right)\left(e^{-j\beta_1(x\sin\theta_i)}\right)\left(-\cos\theta_i\,\hat{x}\right) + \left(\frac{E_{\perp r}}{\eta_1}\right)\left(e^{-j\beta_1(x\sin\theta_r)}\right)\left(\cos\theta_r\,\hat{x}\right) = \left(\frac{E_{\perp t}}{\eta_2}\right)\left(e^{-j\beta_2(x\sin\theta_t)}\right)\left(-\cos\theta_t\,\hat{x}\right)$$

In the equation above, we have used the tangential magnetic field vector epression given in slide 23 Starting equation which is found in slide 24 and repeated below is different from the equation in slide 16

 $(\frac{E_{\perp i}}{n_1})(e^{-j\beta_1(x\sin\theta_i)})(-\cos\theta_i\,\hat{x}) + (\frac{E_{\perp r}}{n_1})(e^{-j\beta_1(x\sin\theta_r)})(\cos\theta_r\,\hat{x}) = (\frac{E_{\perp t}}{n_1})(e^{-j\beta_2(x\sin\theta_t)})(-\cos\theta_t\,\hat{x})$ BUT THE Conlusions ARE THE SAME, AS EXPECTED Conclusion 2) Conclusion 1) $\beta_1(x\sin\theta_i) = \beta_2(x\sin\theta_t)$ $\omega_{\sqrt{\varepsilon_1 \mu_1}}(x \sin \theta_i) = \omega_{\sqrt{\varepsilon_2 \mu_2}}(x \sin \theta_t)$ $\beta_1(x\sin\theta_i) = \beta_1(x\sin\theta_r) = \beta_2(x\sin\theta_t)$ $\omega \sqrt{\varepsilon_{r_1} \varepsilon_0 \mu_{r_1} \mu_0} (x \sin \theta_i) = \omega \sqrt{\varepsilon_{r_2} \varepsilon_0 \mu_{r_2} \mu_0} (x \sin \theta_t)$ $\theta_i = \theta_r$ $\omega n_1(x\sin\theta_i) = \omega n_2(x\sin\theta_t)$ $\omega \sqrt{\varepsilon_{r1} \mu_{r1}} (x \sin \theta_i) = \omega \sqrt{\varepsilon_{r2} \mu_{r2}} (x \sin \theta_t)$ $n_1(\sin\theta_i) = n_2(\sin\theta_t)$ Phase is preserved at the boundary Snell's Law n_1 : refractive index of the 1st medium

ELES 15 Electromegones fractive index of the 2nd medium

$$(\frac{E_{\perp i}}{\eta_1})(e^{-j\beta_1(x\sin\theta_i)})(-\cos\theta_i\,\widehat{x}) + (\frac{E_{\perp r}}{\eta_1})(e^{-j\beta_1(x\sin\theta_r)})(\cos\theta_r\,\widehat{x}) = (\frac{E_{\perp t}}{\eta_2})(e^{-j\beta_2(x\sin\theta_t)})(-\cos\theta_t\,\widehat{x})$$

Using the starting equation found in slide 24 (and rewritten above for convenience), we obtain the second equation which comes from the boundary condition for the tangential components of the magnetic field vectors. This second equation, together with the first equation found in slide 18 ($E_{\perp i} + E_{\perp r} = E_{\perp t}$), will be used to express the unknown quantities $E_{\perp r}$ and $E_{\perp t}$ in terms of the known quantities like $E_{\perp i}$, θ_i and θ_t :

The equation at the top can be simplied as the following using the Conclusion 1 in slide 25

$$\left(\frac{E_{\perp i}}{\eta_1}\right)\left(-\cos\theta_i\,\widehat{x}\right) + \left(\frac{E_{\perp r}}{\eta_1}\right)\left(\cos\theta_r\,\widehat{x}\right) = \left(\frac{E_{\perp t}}{\eta_2}\right)\left(-\cos\theta_t\,\widehat{x}\right)$$

Dropping the \hat{x} on both sides we obtain the equality below which is the 2nd Equation

Again

$$(\frac{E_{\perp i}}{\eta_1})\mathbf{cos}\theta_i - (\frac{E_{\perp r}}{\eta_1})\mathbf{cos}\theta_r = (\frac{E_{\perp t}}{\eta_2})\mathbf{cos}\theta_t$$
a21

New

 $E_{\perp i} + E_{\perp r} = E_{\perp t} \qquad a20$ $\theta_i = \theta_r \text{ (found before)}$ $=> \cos(\theta_i) = \cos(\theta_r) =>$ $\left(\frac{E_{\perp i}}{\eta_1} - \frac{E_{\perp r}}{\eta_1}\right) \cos(\theta_i) = \frac{E_{\perp t}}{\eta_2} \cos(\theta_t) \qquad a21$





New

Check
$$(1 + \Gamma_{\!\scriptscriptstyle \perp}) = \tau_{\!\scriptscriptstyle \perp}$$

Summary, **L** Polarization



From Slide 17 or Slide 24