


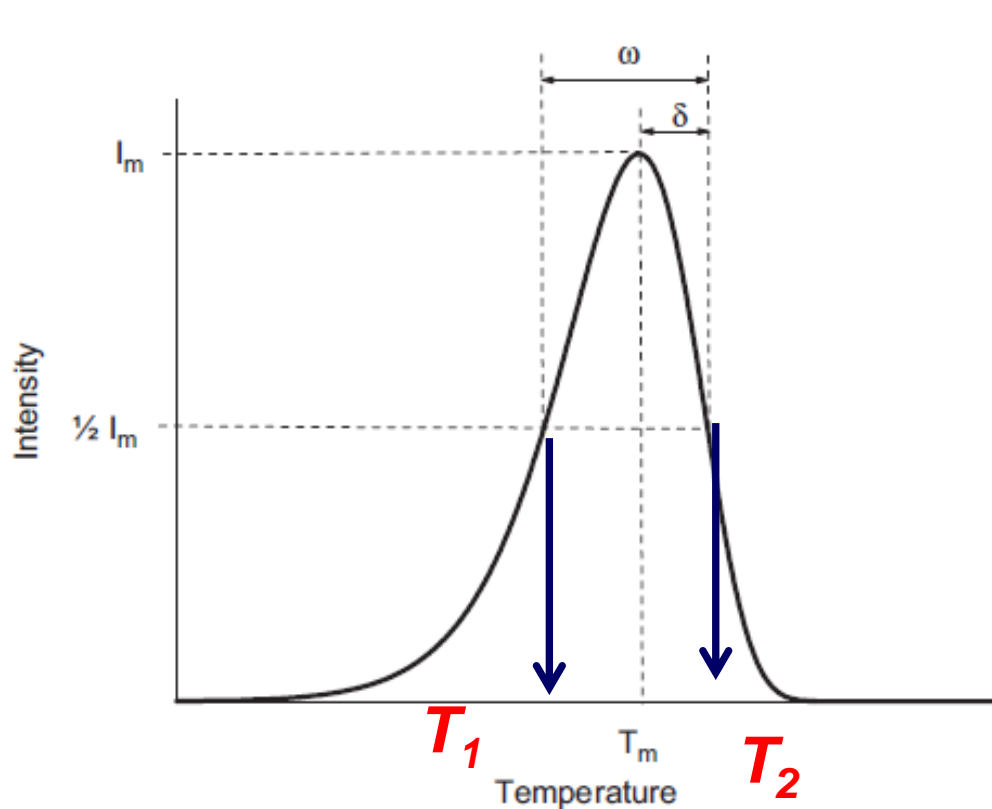
Dependence on the heating rate β

$$I_{TL}(T) = s \cdot \frac{n_0^b p_r}{N p_t} \cdot \exp\left(-\frac{E}{kT}\right) \cdot \left(1 + \frac{p_r s}{p_t N} \cdot \frac{n_0^{b-1}}{\beta} \cdot (b-1) \cdot \int_{T_0}^T \exp\left(-\frac{E}{kT'}\right) dT'\right)^{\frac{b}{1-b}}$$


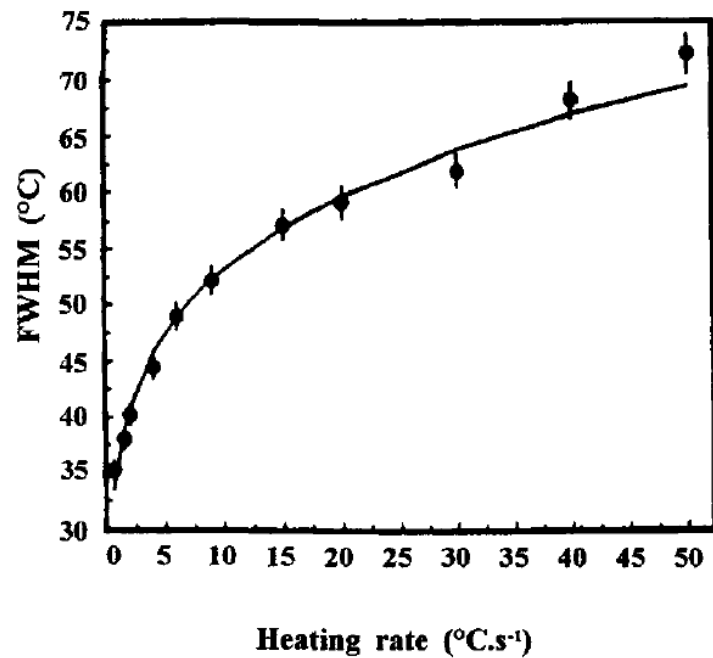
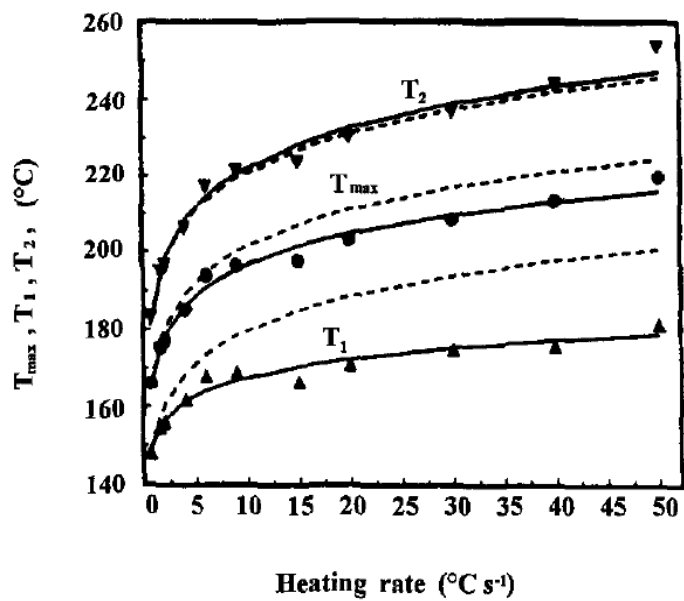
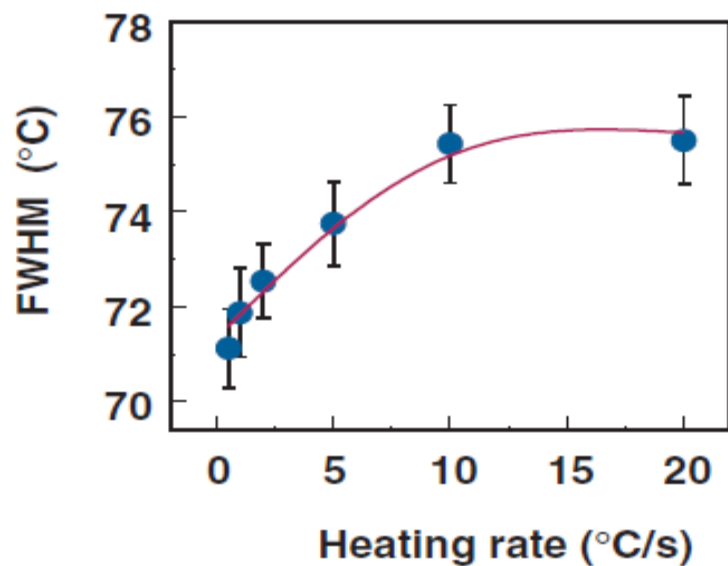
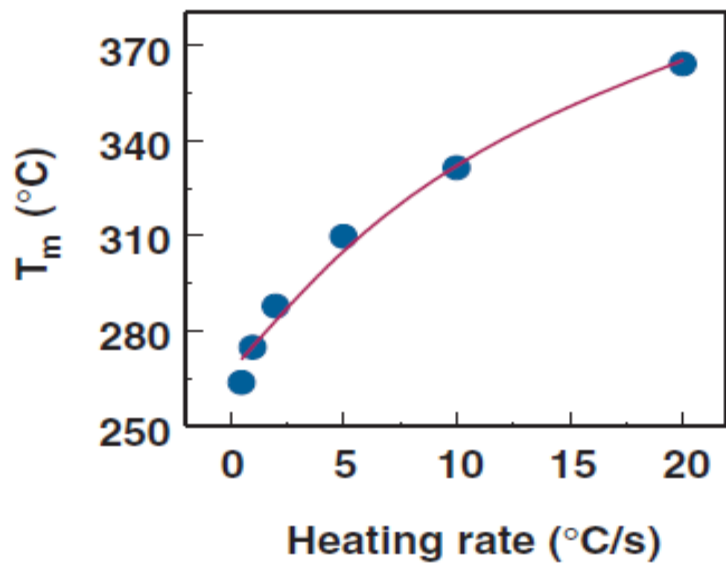
Normalisation factor so that the following integral equals to 1, since it represents the probability that all electrons are de-trapped

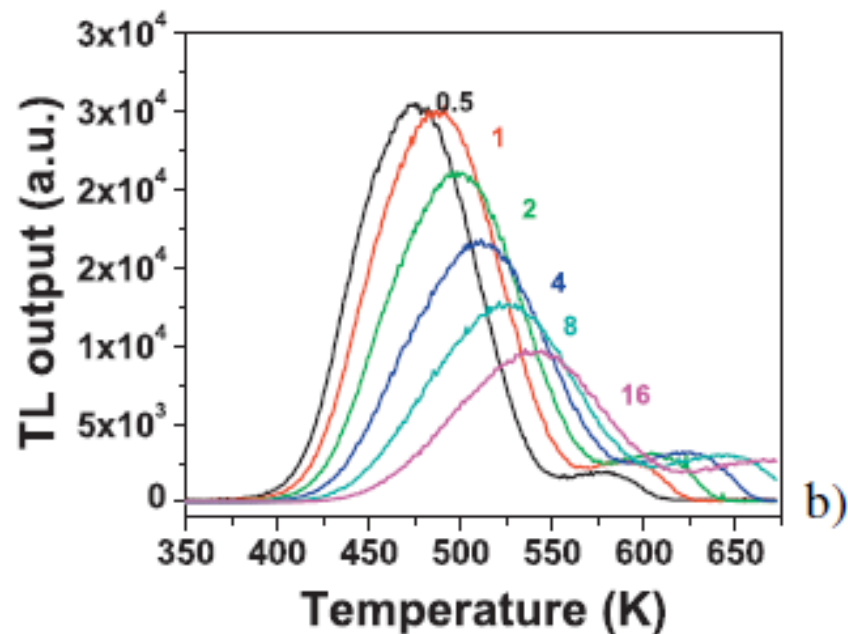
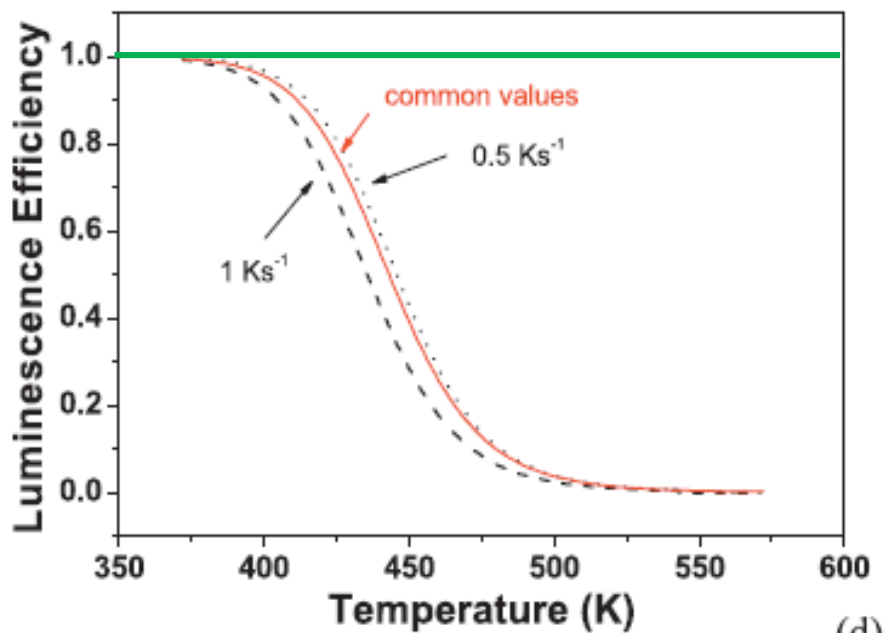
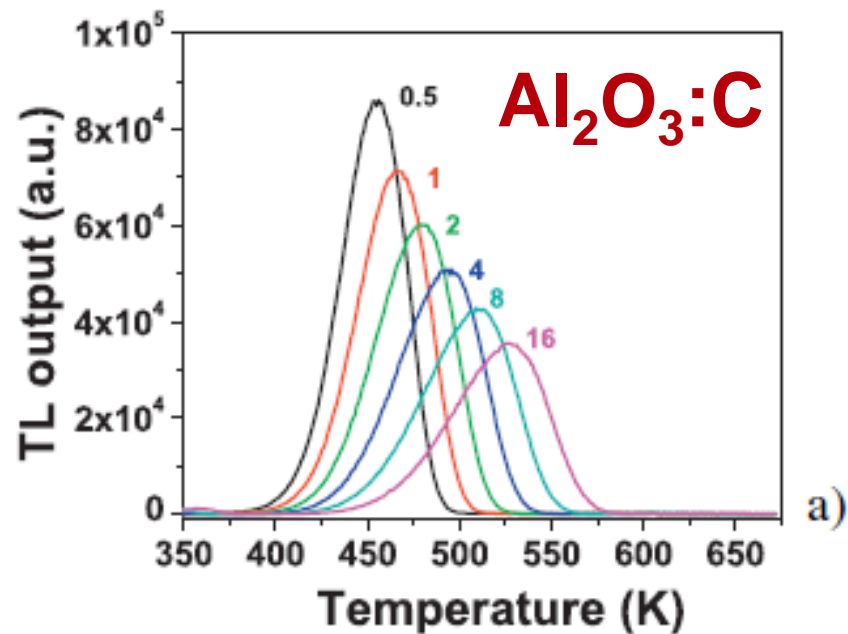
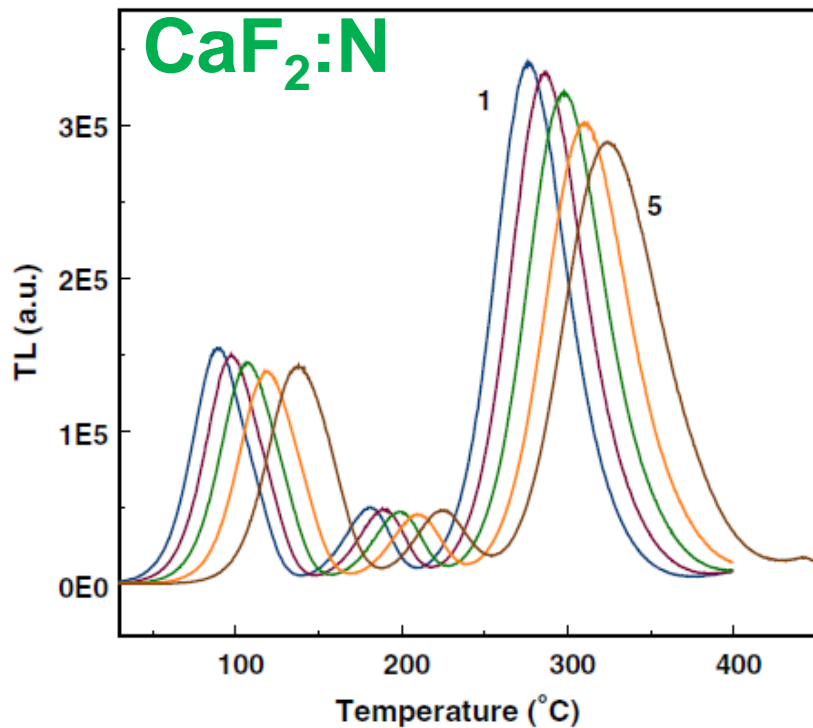
The H.R. effect on TL glow-peaks has been largely discussed by G.Kitis [3] who considers the H.R. as a **dynamic parameter** rather than a simple experimental setup variable. His study has been carried out on single, well separated glow peaks, considering the following experimental characteristics: i.e., T_M , full width at half maximum (FWHM), peak intensity and peak integral.

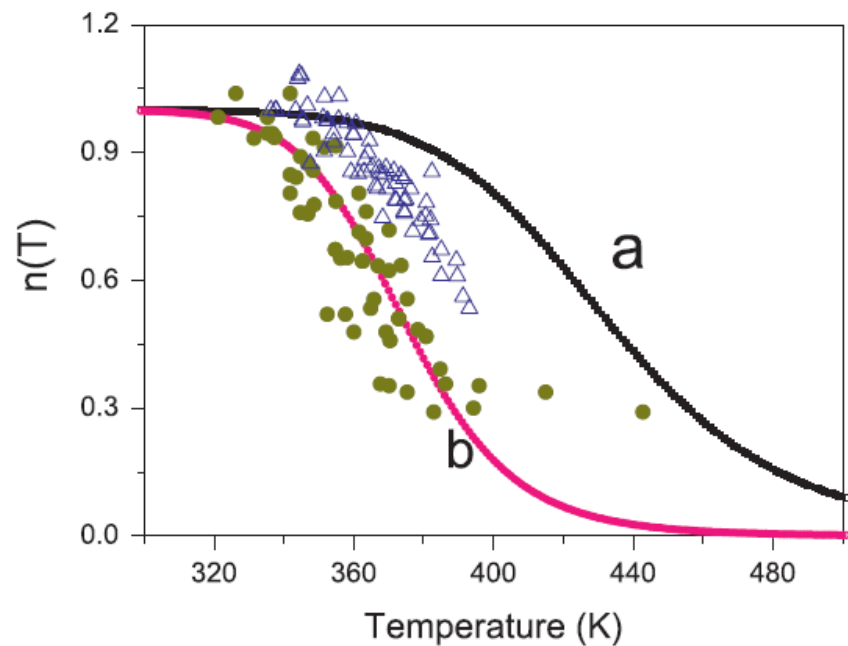
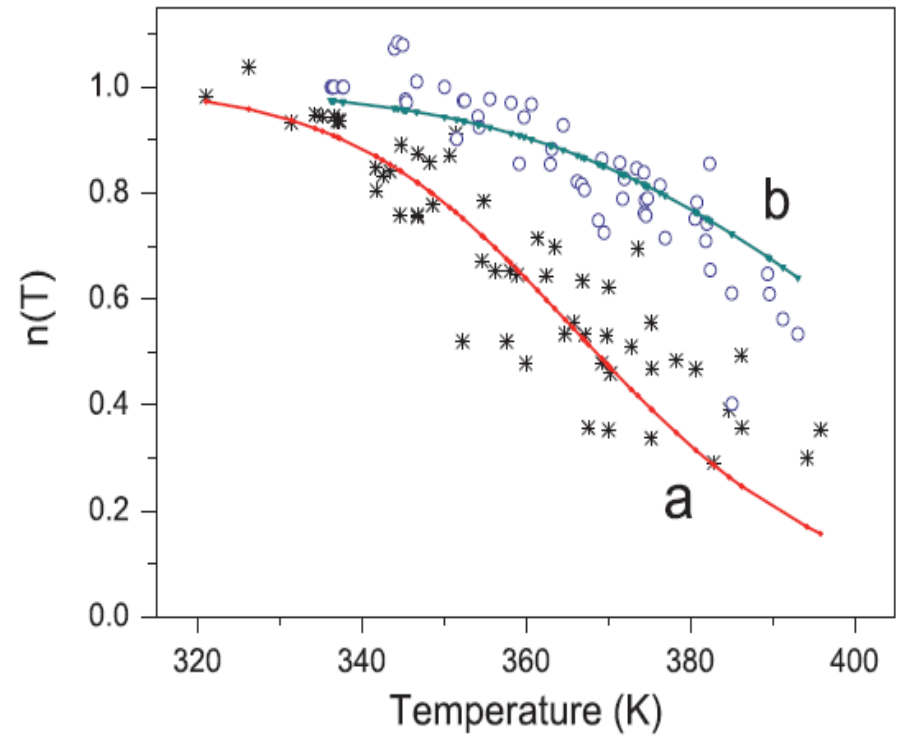
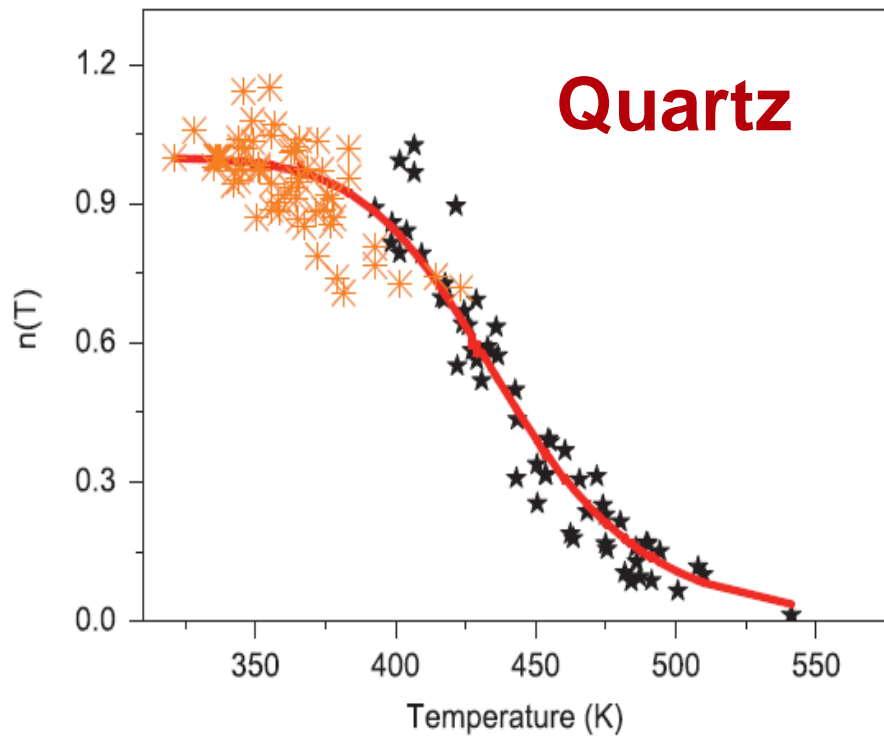
- *The maximum glow-peak temperature T_m is shifted towards higher temperatures as the heating rate increases.*
- *The same takes place for T_1 and T_2 . Therefore, the FWHM increases for increasing heating rate.*



- *The integrated intensity of each single peak stays constant. → Otherwise thermal quenching*

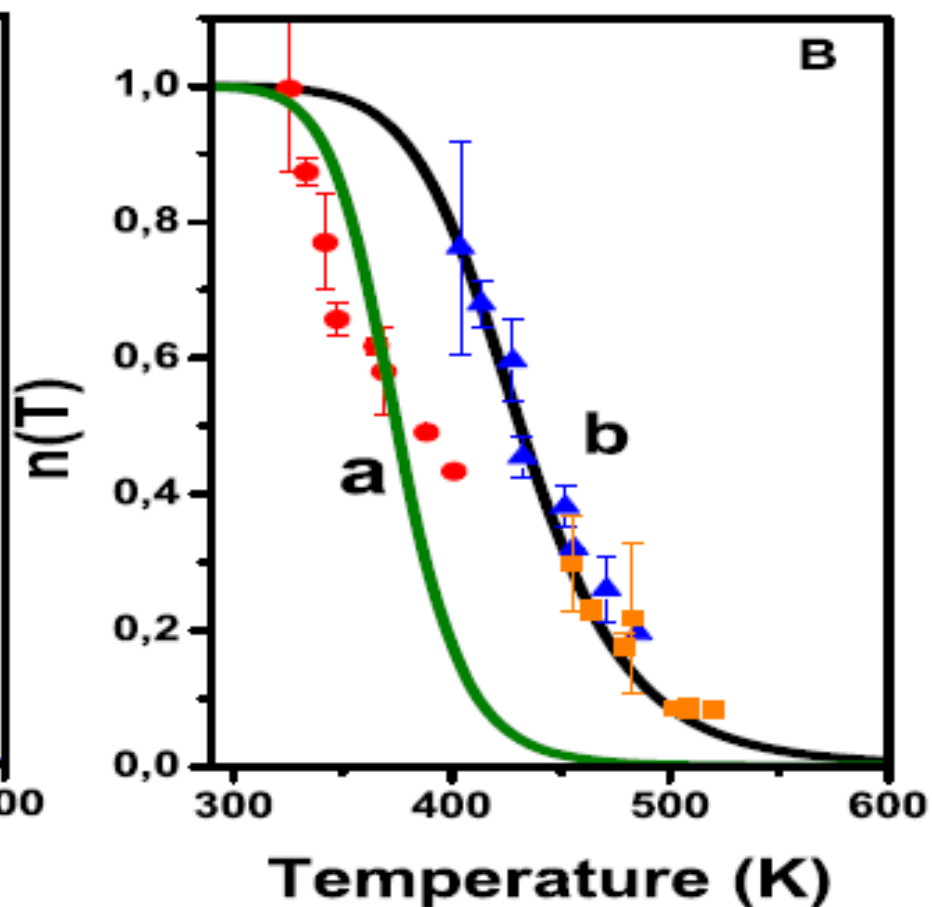
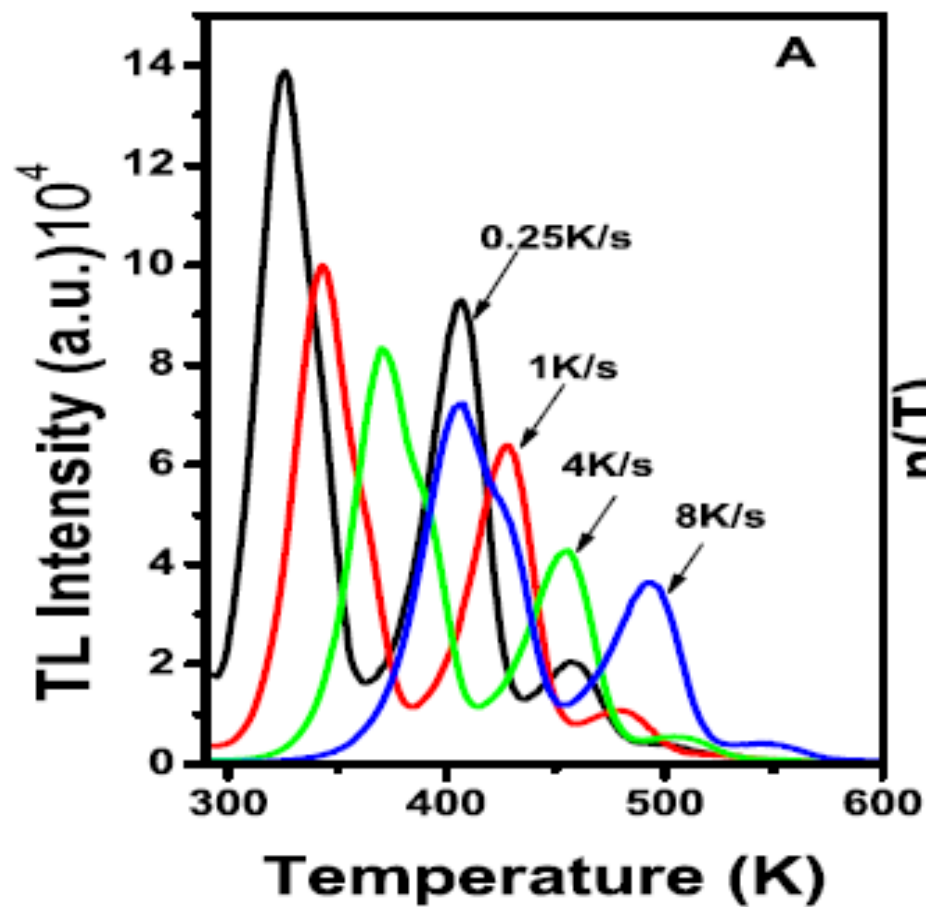






$$\eta = \frac{P_r}{P_r + P_{nr}}$$

$$n = \frac{1}{1 + c \exp\left(-\frac{\Delta E}{kT}\right)}$$



Quartz

Various heating rates β method

$$\frac{\beta_1 E}{kT_{M1}^2} = s \exp\left(-\frac{E}{kT_{M1}}\right)$$

$$\frac{\beta_2 E}{kT_{M2}^2} = s \exp\left(-\frac{E}{kT_{M2}}\right)$$

$$E = k \frac{T_{M1} T_{M2}}{T_{M1} - T_{M2}} \ln\left(\frac{\beta_1}{\beta_2}\right) \cdot \left(\frac{T_{M2}}{T_{M1}}\right)^2$$

First order

$$\left(\frac{\beta}{T_M^2}\right) \cong \exp\left(-\frac{E}{kT_M}\right) \left(\frac{ks}{E}\right) [1 + (b-1)\Delta_M]$$

General order