

# ENE 505 – Applied Computational Fluid Dynamics in Renewable Energy Technologies

## WEEK 8: TURBULENCE MODELS

### TURBULENCE MODELS:

- **RANS based Turbulence Models**

- For turbulent, incompressible and Newtonian viscous fluid, the Reynolds-averaged Navier–Stokes equation is given by:

$$\rho \frac{D\bar{u}_i}{Dt} = -\frac{\partial \bar{p}}{\partial x_i} + \frac{\partial}{\partial x_j} \left[ \mu \left( \frac{\partial \bar{u}_i}{\partial x_j} + \frac{\partial \bar{u}_j}{\partial x_i} - \frac{2}{3} \delta_{ij} \frac{\partial \bar{u}_i}{\partial x_i} \right) \right] + \frac{\partial}{\partial x_j} \left( -\rho \overline{u_i' u_j'} \right)$$

Where  $\rho$  is the fluid density,  $\bar{u}_i$  is the time averaged velocity,  $u_i'$  is the deviation from the time averaged velocity,  $\bar{p}$  is the time averaged pressure,  $\mu$  is the dynamic viscosity of the fluid,  $-\rho \overline{u_i' u_j'}$  is the Reynold's stress tensor, which appears on the right-hand side of the time-averaged Navier–Stokes equations as a result of time averaging to Navier–stokes equations.

- The temporal and spatial co-ordinates correspond to  $t$  and  $x_j$ , respectively. In an eddy-viscosity model, the Reynolds stress is assumed to be proportional to the mean velocity gradients, with the constant of proportionality being the turbulent viscosity,  $\mu_t$ . This assumption is known as the Boussinesq eddy-viscosity hypothesis, and provides the following expression for the Reynolds stresses [1]

$$-\rho \overline{u_i' u_j'} = \frac{2}{3} \rho k \delta_{ij} + \mu_t \left( \frac{\partial \bar{u}_i}{\partial x_j} + \frac{\partial \bar{u}_j}{\partial x_i} \right) + \frac{2}{3} \mu_t \frac{\partial \bar{u}_j}{\partial x_j} \delta_{ij}$$

where  $k$  is the turbulent kinetic energy expressed as:

$$k = \frac{1}{2} \overline{u_i u_j}$$

- **Standard k-epsilon Model**

- The standard k- $\epsilon$  turbulence model is based on transport equations for turbulent kinetic energy  $k$  and its dissipation rate,  $\epsilon$ , and it is firstly developed by Launder and Spalding [2].
- The turbulence kinetic energy,  $k$  and its rate of dissipation,  $\epsilon$  are obtained from the following transport equations due to [2]

$$\rho \frac{Dk}{Dt} = \frac{\partial}{\partial x_i} \left[ \left( \mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_i} \right] + G_k - \rho \epsilon$$

$$\rho \frac{D\epsilon}{Dt} = \frac{\partial}{\partial x_i} \left[ \left( \mu + \frac{\mu_t}{\sigma_\epsilon} \right) \frac{\partial \epsilon}{\partial x_i} \right] + C_{1\epsilon} \frac{\epsilon}{k} G_k - C_{2\epsilon} \rho \frac{\epsilon^2}{k}$$

where,  $\mu_t$  is the eddy viscosity,  $C_{1\epsilon}$  and  $C_{2\epsilon}$  are constants in the sense that they are not changed between calculations.  $\sigma_k$  and  $\sigma_\epsilon$  are the turbulent Prandtl numbers for  $k$  and  $\epsilon$ , respectively. The default values for these coefficients are given as follows:  $C_{1\epsilon} = 1.44$ ,  $C_{2\epsilon} = 1.92$ ,  $C_\mu = 0.09$ ,  $\sigma_k = 1.0$ , and  $\sigma_\epsilon = 1.0$ .  $G_k$  represents the generation of turbulent kinetic energy due to the mean velocity gradients, calculated in a manner consistent with the Boussinesq hypothesis.

- **The RNG k-epsilon Model**

- Further development and Improvement from the standard model has been done by Yakhot et al. [3] and is based on the renormalized (RNG) group theory. An additional sink term is suggested; over the standard k- $\epsilon$  turbulence model, in the turbulence dissipation equation to account for non-equilibrium strain rates and employs different values for various model coefficients.

- The equation of  $k$  remains the same and the dissipation equation  $\varepsilon$  is modified to include the additional sink term:

$$\frac{C_\mu \eta^3 (1 - \eta/\eta_0) \varepsilon^3}{1 + \beta \eta^3} k$$

- **The k-omega Model**

- The  $k$ - $\omega$  model is based on Wilcox [4]. It compromises modifications for low Reynolds number effects, compressibility and shear flow spreading. It is characterized by the turbulent kinetic energy and the frequency,  $\omega = k/\varepsilon$ , where  $\varepsilon$  is the rate of dissipation of  $k$ . The turbulence viscosity can be expressed as:

$$\mu_t = C_\mu \frac{\rho k}{\omega}$$

where the turbulent kinetic energy,  $k$ , and the specific dissipation rate,  $\omega$ , are obtained from the following transport equations:

$$\frac{\partial \rho k}{\partial t} + \frac{\partial (\rho \bar{u}_i k)}{\partial x_j} = \frac{\partial}{\partial x_j} \left[ \left( \mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] + P - \rho \omega k$$

$$\frac{\partial \rho \omega}{\partial t} + \frac{\partial (\rho \bar{u}_i \omega)}{\partial x_j} = \frac{\partial}{\partial x_j} \left[ \left( \mu + \frac{\mu_t}{\sigma_\omega} \right) \frac{\partial \omega}{\partial x_j} \right] + \gamma \frac{\omega}{k} - \beta \rho \frac{\omega^2}{C_\mu}$$

**References:**

1. Hinze, O. (1975) Turbulence. New York, McGraw-Hill Publishing Co.
2. Launder BE, Spalding DB. Mathematical Models of Turbulence. Academic Press: London, 1972.
3. Yakhot V, Orszag SA, Thangham S, Gatski TB, Speziale CG. Development of turbulence models for shear flows

by a double expansion technique. *Physics of Fluids* 1992; 4(7): 1510–1520.

4. Wilcox, D. (1998), *Turbulence modeling for CFD*, California: DCW Industries, Inc., La Canada.